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FEEDBACK STABILIZATION OF NAVIER–STOKES EQUATIONS

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Abstract. The study of the local exponential stabilization problem for the Navier-Stokes equations using the algebraic Riccati equation is the main aim of this paper.
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1. INTRODUCTION

Consider the Navier–Stokes system

(1.1)

$$y_t(x,t) - \nu \Delta y(x,t) + (y \cdot \nabla) y(x,t) = m(x)u(x,t) + f_0(x) + \nabla p(x,t), \quad (x,t) \in Q$$

$$(\nabla \cdot y)(x,t) = 0, \qquad \forall (x,t) \in Q = \Omega \times (0,\infty)$$

$$y = 0, \qquad \text{on } \Sigma = \partial \Omega \times (0,\infty)$$

$$y(x,0) = y_0(x), \qquad x \in \Omega$$

in a domain Ω of R^d , d = 2 or d = 3 where $y = (y_1, ..., y_d)$ is the state $u = (u_1, u_2, ..., u_d)$ is the control input, p is the unknown pressure and m is the characteristic function of an open subset ω of Ω . Let

$$\begin{split} H &= \{y \in (L^2(\Omega))^d, \ \nabla \cdot y = 0, \ y \cdot n = 0 \text{ on } \partial \Omega \}, \ V = \{y \in (H^1_0(\Omega))^d, \ \nabla \cdot y = 0 \} \\ \text{and } P : (L^2(\Omega))^d \longrightarrow H \text{ be the Leray projector.We set} \end{split}$$

$$b(y,z,w) = \sum_{i,j=1}^{d} \int_{\Omega} y_i D_i z_j w_j dx, \ (By,w) = b(y,y,w), \ \forall y,w \in V$$

and rewrite equation (1.1) as

(1.1)'
$$\frac{dy}{dt}(t) + \nu Ay(t) + By(t) = P(mu) + Pf_0, \quad t \in [0, \infty)$$

$$y(0) = y_0, \qquad \qquad t \in (0, \infty)$$

where $A \in L(V, V')$ (the Stokes operator)

$$(Ay, w) = \sum_{i=1}^{d} \int_{\Omega} \nabla y_i \cdot \nabla w_i dx, \ \forall y, w \in V.$$

Let (y_e, p_e) be a steady-state solution to (1.1), i.e.,

$$\begin{aligned} -\nu \Delta y_e + y_e \cdot \nabla y_e &= \nabla p_e + f_0(x) & \text{in } \Omega \\ \nabla \cdot y_e &= 0 & \text{in } \Omega \\ y_e &= 0 & \text{on } \partial \Omega. \end{aligned}$$

We assume that the boundary $\partial \Omega$ is a finite union of d-1 dimensional C^{∞} connected manifolds diffeomorphic with $S_d^r = \{x \in \mathbb{R}^d, |x| = r\}.$

For d = 2, 3 always there is at least one steady-state solution but for small viscosity constant ν this solution might be instable. However, by a recent result of O.Imanuvilov [6], [7] (see also [2]) if (y_e, p_e) and y_0 are sufficiently smooth, for instance if

(1.2)
$$(y_e, p_e) \in ((H^3(\Omega))^d) \cap V) \times H^1(\Omega), \ y_0 \in (H^2(\Omega))^d \cap V$$

and if $||y_0 - y_e||_{(H^2(\Omega))^d} \leq \eta$ is sufficiently small then for each T > 0 there are

(1.3)
$$u \in H^1(0,T; (L^2(\Omega))^d), y \in L^\infty(0,T; (H^2(\Omega))^d \cap V) \cap H^1(0,T; H)$$

and $p \in L^2(0, T; H^1(\Omega))$ satisfying (1.1) and such that $y(x, T) \equiv y_e(x)$. In particular, this implies that there is a controller u which stabilizes the steady-state solution y_e . Here we shall treat the local exponential stabilization problem for the Navier–Stokes using the algebraic Riccati equation associated with the linearized Stokes equation.T he argument will be scketched only; the complete proofs will appear elsewhere. For recent results on stabilization of fluid flows we refer to the works [1],[3],[4].

Here $H^k(\Omega)$ and $H^1(0,T;X)$, (X is a Hilbert space) are usual Sobolev spaces. We denote by $|\cdot|$ the norm of H and $(L^2(\Omega))^d$ and by $||\cdot||$ the norm of V. By (\cdot, \cdot) denote the pairing between V, V' (the dual space of V) and, respectively, the scalar product of H. Finally, $|\cdot|_s$ is the norm of the Sobolev space $(H^s(\Omega))^d$.

2. STABILIZATION OF THE LINEARIZED EQUATION

Substituting y by $y + y_e$ into (1.1) we are lead to the equation

(2.1)
$$\begin{array}{l} y_t - \nu \Delta y + (y \cdot \nabla)y + (y_e \cdot \nabla)y + (y \cdot \nabla)y_e = mu + \nabla p & \text{in } Q \\ \nabla \cdot y = 0 & \text{in } Q \\ y = 0 & \text{on } \Sigma \\ y(x,0) = y_0(x) - y_e(x) = y^0(x), & x \in \Omega. \end{array}$$

(2.2)
$$\frac{dy}{dt}(t) + \nu Ay(t) + A_0y(t) + By(t) = P(mu), \ t \ge 0$$
$$y(0) = y_0$$

where $A_0 \in L(V, H)$ is defined by

(2.3)
$$(A_0y, w) = b(y_e, y, w) + b(y, y_e, w), \ \forall w \in H$$

Consider the linear system

(2.4)
$$\frac{dy}{dt}(t) + \nu Ay(t) + A_0 y(t) = P(mu)(t), \ t \ge 0, \ y(0) = y^0$$

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and the optimal control problem

(2.5)
$$\varphi(y^0) = \operatorname{Min}\left\{\frac{1}{2}\int_0^\infty \left(\left|A^{\frac{7}{8}}y(t)\right|^2 + \left|u(t)\right|^2\right)dt; \ (y,u) \text{ subject to } (2.4)\right\}.$$

Denote by $D(\varphi)$ the set of all $y^0 \in H$ for which $\varphi(y^0) < \infty$. Under our assumptions for each $y^0 \in H$ the linear Stokes equation (2.4) is exactly null controllable (see [6], [7].) More precisely, there are $u \in H^1(0,T; (L^2(\Omega))^d), y \in L^2(0,T; (H^2(\Omega))^d \cap V)$ solution to (2.4) such that $y(T) \equiv 0$ and (see Lemma 3.1 in [2])

$$|u|_{H^1(0,T;(L^2(\Omega))^d)} \le C |y^0|.$$

Moreover, we have

(2.6)
$$\varphi(y^0) \ge C \left| A^{\frac{3}{8}} y^0 \right|^2.$$

because (see e.g. [5], [8])

$$(A_0 y, A^{\frac{3}{4}} y) \bigg| \le \bigg| b(y, y_e, A^{\frac{3}{4}} y) \bigg| + \bigg| b(y_e, y, A^{\frac{3}{4}} y) \bigg| \le C ||y|| \bigg| A^{\frac{3}{4}} y \bigg|.$$

By (2.6) we see that $D(\varphi) = D(A^{\frac{3}{8}}) = W$. The space W is endowed with the graph norm $|y|_W = |A^{\frac{3}{8}}y|$. Thus there is a linear self-adjoint operator $R: H \to H$ with the domain $D(R) = D(A^{\frac{3}{4}})$ such that

$$\frac{1}{2}(Ry^0, y^0) = \varphi(y^0), \ \forall \, y^0 \in D(A^{\frac{3}{4}}).$$

Moreover, $R \in L(W, W')$ and the latter extends to all of W.

PROPOSITION 1. Let d = 2, 3. Then the optimal control u^* is expressed as (2.7) $u^*(t) = -mRy^*(t), \ \forall t > 0.$

Moreover, there is $\omega > 0$ such that

(2.8)
$$(Ry, y) \ge \omega \|y\|^2_W, \ \forall y \in W$$

and R is the solution to algebraic Riccati equation

(2.9)
$$(\nu Ay + A_0y, Ry) + \frac{1}{2} |mRy|^2 = \frac{1}{2} |A^{\frac{7}{8}}y|^2, \ \forall y \in D(A).$$

Proof. Estimate (2.8) follows by (2.6). For T > 0, let (y^*, u^*) be the solution to optimal control problem

$$\operatorname{Min}\left\{\frac{1}{2}\int_{0}^{T} \left(\left|A^{\frac{7}{8}}y(t)\right|^{2} + \left|u(t)\right|^{2}\right) dt + \varphi(y(T)); (y, u) \text{ subject to } (2.4) \text{ on } (0, T)\right\}.$$

Thus $u^*(t) = mq^T(t), \ \forall t \in [0,T)$ where $q^T \in L^2(0,T;H) \cap C_w([0,T];V')$ is the solution to adjoint equation

(2.10)
$$\frac{d}{dt}q^{T} - (\nu A + A_{0})^{*}q^{T} = A^{\frac{7}{4}}y^{*}, \ t \in (0,T)$$
$$q^{T}(T) = -Ry^{*}(T).$$

By (2.10) and the unique continuous property for the Stokes equation it follows that $q^T = q^{T'}$ on (0,T) for 0 < T < T'. Hence $q^T = q$ is independent of T and so the above equations extend to all of R^+ . Moreover,

$$q(t) = -Ry^*(t), \ \forall t \ge 0$$

and so we obtain (2.7) as claimed. Next we have

$$\begin{split} \varphi(y^*(t)) &= \frac{1}{2} \int_t^\infty (\left|A^{\frac{7}{8}}y^*\right|^2 + |u^*|^2) ds, \ \forall t \ge 0\\ \left(Ry^*(t), \frac{dy^*}{dt}(t)\right) &+ \frac{1}{2} \left|A^{\frac{7}{8}}y^*(t)\right|^2 + \frac{1}{2} |mRy^*(t)|^2 = 0, \ \text{a.e.} \ t > 0. \end{split}$$

This yields

$$-(Ry^{*}(t),\nu Ay^{*}(t)+A_{0}y^{*}(t))-\frac{1}{2}|mRy^{*}(t)|^{2}+\frac{1}{2}|A^{\frac{7}{8}}y^{*}(t)|^{2}=0, \ \forall t\geq 0.$$

3. STABILIZATION OF THE NAVIER-STOKES EQUATION

THEOREM 1. Let d = 2,3 and let R be the operator defined in Proposition 1. Let $(y_e, p_e) ((H^3(\Omega))^d \cap V) \times H^1(\Omega)$ be a steady-state solution to equation (1.1). Then the feedback controller

$$(3.1) u = -mR(y - y_e)$$

exponentially stabilizes y_e in a neighbourhood $\mathcal{V} = \{y_0 \in W; \|y_0 - y_e\|_W < \rho\}$ of y_e . More precisely, for each $y_0 \in \mathcal{V}$ there is a weak solution $y \in L^{\infty}_{\text{loc}}(R^+; H) \cap L^2_{\text{loc}}(R^+; V)$ to closed loop system

(3.2)
$$\frac{dy}{dt} + \nu Ay + A_0 y + By + P(mR(y - y_e)) = Pf_0, \ t > 0$$
$$y(0) = y_0$$

such that

(3.3)
$$|y(t) - y_e| \le C e^{-\gamma t} ||y_0 - y_e||_W, \ \forall y_0 \in \mathcal{V}$$

for some $\gamma > 0$.

Proof. We reduce the problem to that of stability of the null solution to corresponding closed loop system (2.2), i.e.,

(3.4)
$$\frac{dy}{dt} + \nu_0 Ay + A_0 y + By + P(mRy) = 0, \ t > 0, \ y(0) = y^0.$$

We consider the approximating equation

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(3.5)
$$\frac{dy_N}{dt} + \nu_0 Ay_N + A_0 y_N + B_N y_N + P(mRy_N) = 0, y_N(0) = y^0,$$

where

$$B_N y = By$$
 if $||y|| \le N$, $B_N y = \frac{N^2}{||y||^2} By$ if $||y|| > N$.

For each $y^0 \in D(A)$ equation (3.5) has a unique solution $y_N \in W^{1,\infty}_{\text{loc}}(R^+;H) \cap L^{\infty}_{\text{loc}}(R^+;D(A))$. If $y^0 \in V$ then

(3.6)
$$|y_N(t)|^2 + \int_0^t ||y_N(s)||^2 ds \le C_T, \ \forall t \in (0,T),$$

where C_T is independent of N. Thus on a subsequence $N \to \infty$

(3.7)
$$y_N \longrightarrow y$$
 weak star in $L^{\infty}_{\text{loc}}(R^+; H) \cap L^2_{\text{loc}}(R^+; V)$

where $y \in L^2_{loc}(R^+; V) \cap C_w(R^+; H)$ is a weak solution to equation (3.4). Now we multiply equation (3.5) by Ry_N and use (2.9) to obtain that

(3.8)
$$\frac{d}{dt}(Ry_N(t), y_N(t)) + |mRy_N(t)|^2 + |A^{\frac{3}{4}}y_N(t)|^2 = 2(B_Ny_N(t), Ry_N(t)), \text{ a.e. } t > 0.$$

On the other hand, we have by (2.7) and the properties of b (see e.g. [5], [8]) that

(3.9)
$$\begin{aligned} |(B_N y_N, y_N)| &\leq \inf\left(1, \frac{N^2}{\|y_N\|^2}\right) |b(y_N, y_N, Ry_N)| \leq \\ &\leq C|y_N|_{\frac{3}{4}}|y_N|_{\frac{7}{4}}|Ry_N| \leq C \left|A^{\frac{3}{8}}y_N\right| \left|A^{\frac{7}{8}}y_N\right| \left|A^{\frac{3}{4}}y_N\right| \leq \\ &\leq C \left|A^{\frac{7}{8}}y_N\right|^2 (Ry_N, y_N)^{\frac{1}{2}} \end{aligned}$$

because $D(A^m) \subset (H^{2m}(\Omega))^n$ for all noninteger m. We set

$$E = \{y^0 \in W; \ (Ry^0, y^0) < \rho\}$$

By (3.8) and (3.9) we see that for ρ sufficiently small and $y^0 \in E$ we have

$$\frac{d}{dt}(Ry_N(t), y_N(t)) + \frac{1}{2} \left| A^{\frac{7}{8}} y_N(t) \right|^2 \le 0, \text{ a.e. } t > 0$$

and this yields

$$|y_N(t)| \le ||y_N(t)||_W \le C ||y^0||_W e^{-\gamma t}, \ \forall t \ge 0$$

for some $\gamma > 0$. Then we see by (3.7) that

$$|y(t)| \le C \left\| y^0 \right\|_W e^{-\gamma t}. \ \forall t \ge 0$$

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